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# 1 **Morphological and force spectroscopy characterizations for indentifi-** 2 **cation of surface nanobubbles from nanodroplets and blisters**

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## 12 **Abstract**

13 Surface nanobubbles (NBs) play an important role in various practical applications, such as min-  
14 eral flotation and separation, drag reduction, and nanostructured surface fabrication. Until now, it  
15 still remains as a challenge to identify surface NBs from other spherical-cap-liked nano-objects,  
16 like blisters and nanodroplets (NDs). Here we focus on the distinctions of NBs from NDs and blis-  
17 ters using an atomic force microscopy. It is implemented through morphological characterization,  
18 high load scanning, and force spectroscopy measurement. In the morphological characterization  
19 experiment, contact angles of the three types of nano-objects were compared. In the high load  
20 scanning experiment, the response of the nano-objects to high scanning loads was studied. The  
21 mobility, deformability, and volume change of the nano-objects during the high load scanning were  
22 investigated. At last, the force spectroscopy measurement was implemented. Due to the existence  
23 of the three-phase contact lines on both tip-NB and tip-ND interactions, force-distance curves ex-

hibit the similar behaviors on both NBs and NDs. However, quantitative analysis shows that the extracted parameters from force-distance curves can be used to distinguish one from the other. This study developed a systematic way to distinguish surface NBs from others nano-objects, which is crucial for surface nanobubble community.

## Keywords

nanobubbles; blisters; nanodroplets; morphological characterization; force spectroscopy measurement; atomic force microscopy

## Introduction

Surface nanobubbles (NBs) at solid liquid interfaces have attracted significant attentions in the last two decades because of their great potential in numerous applications, such as mineral flotation and separation [1], drag reduction [2-4], nanostructured surface fabrication [5-8] and the context of catalysis and electrolysis [9,10]. The properties of NBs have been investigated with numerous techniques, including atomic force microscopy (AFM) [11-13], x-ray reflectivity [14,15], infrared spectroscopy [16,17], and optical microscopy [18-20]. However, the NB community has long been suffering from the confusion caused by some other spherical cap shaped nano-objects, like nanodroplets (NDs) [21,22] and blisters [23]. NDs may nucleate in NB experiments because of the impurities in liquid. This is because NDs basically are similar to NBs regarding the nucleation mechanism. The blisters are thin polymer film wrapped water pockets at solid-liquid interfaces on a thin film coated silicon substrate [24,25]. They are produced as water permeates through the thin film (i.e. a polystyrene film), wet and hence detach the thin films, leading to the formation of a water reservoir in between the supporting silicon substrate and the thin film [23,26]. So far it is still difficult to distinguish NBs from these spherical cap shaped nano-objects, especially from NDs. NDs and NBs are all soft in nature, which makes it more challenging to distinguish one from the other with current imaging techniques.

Several research works have been performed to identify surface NBs [22,27-30]. The gaseous NBs

49 are expected to dissolve over time in degassed water. Thus the response of the nucleated spherical  
50 objects in degassed water was used to confirm the existence of surface nanobubbles [22,27]. The  
51 results show that NBs indeed could disappear in degassed water. However, Zhang *et al.* found that  
52 NBs can survive even after several hours of degassing [27]. This may be due to the cooperative  
53 shielding effects from neighboring NBs [28].

54 In Chan *et al.*'s study, they compared the interaction of retracting contact lines with NBs, NDs, and  
55 nanoparticles, the dynamic processes of which were recorded by a total internal reflection fluores-  
56 cence (TIRF) microscopy [29]. Considerable different behaviors among NBs, NDs and nanopar-  
57 ticles were observed. Surface NBs rapidly collapsed while contacting with the retracting contact  
58 lines, while an accelerated receding and a pinning-depinning process were observed on the NDs  
59 and NPs, respectively. Meanwhile, Seo et al [30]. distinguished NBs from oil NDs covered by dyes  
60 with a fluorescence microscopy. The gas NBs and oil NDs absorb different fluorescent dyes, emit-  
61 ting fluorescence signals. All these methods provide evidence on the gas nature of NBs, with a lim-  
62 ited spatial resolution though. However, the addition of dyes unavoidably alters the physicochemical  
63 properties of NBs or even causes their collapse, which is undesired in most of systems.

64 Recently, efforts were put on the distinction of surface NBs from other spherical cap liked ob-  
65 jects by using AFMs, which guarantees high spatial resolution [31,32]. Wang *et al* produced poly-  
66 dimethylsiloxane (PDMS) NDs on highly oriented pyrolytic graphite (HOPG) [32]. By comparing  
67 their properties with that of NBs using an AFM, they found that PDMS NDs were very close to  
68 NBs in size, contact angle, and stiffness, but different from force-distance curves. They claimed  
69 that there was a plateau in the approach force distance curves acquired on NBs, but not on the  
70 PDMS NDs. In the tip-nanodroplet interaction process, a linear force-depth curve was observed.

71 An *et al.* also identified NBs from PDMS NDs with an AFM [31]. Their results show that NBs and  
72 NDs exhibit distinct responses under higher vertical and lateral forces. In PeakForce mode, they  
73 found that NBs are invisible in AFM height images under large vertical imaging forces in the or-  
74 der of tens of nN, but NDs maintain a nanometric molecular layer. They also found that NBs are  
75 strongly pinned on substrates and survive large lateral imaging forces of up to 50 nN without being

76 moved or destroyed in the contact mode AFM. NDs are weakly pinned and can be easily moved.  
77 However, in Wang *et al.*'s work, they observed a linear dependence of AFM cantilever deflection,  
78 which is proportional to force, on the penetration depth of AFM tips into NBs. While the interac-  
79 tion curve on nanodroplet shows sharp kinks. This is in conflict with what was reported in An's  
80 work.

81 In this work, the spontaneous formation of NBs on a hydrophobic surface was adopted to avoid  
82 addition of any other solvent and then minimize the chance of contamination. By employing nano-  
83 manipulation and force spectroscopy measurement of AFM tips on NBs, NDs, and blisters, a sys-  
84 tematic investigation was conducted to distinguish NBs from NDs and blisters. We find that NBs  
85 are distinguishable from three aspects: (a) volume changes before and after coalescence, (b) re-  
86 sponse to higher loads, and (c) force spectroscopy measurement.

## 87 **Experimental section**

### 88 **(i) Sample preparation**

89 The polystyrene (PS) surfaces used for NB and blister nucleation were prepared by spin coating  
90 thin films of PS on silicon (100) substrates at a speed of 1000 rpm. Before spin coating, the sub-  
91 strates were sequentially cleaned in sonication bathes of piranha, acetone, and then water, each  
92 for 30 min. PS particles (molecular weight 350 000, Sigma-Aldrich) were dissolved in toluene  
93 (Mallinckrodt Chemical) to make the PS solutions. Two PS films, PS sample 1 and PS sample 2,  
94 were prepared. The PS concentrations for sample 1 and sample 2 are 1.0% and 0.5%, respectively.  
95 The lower PS concentrations leads to a thinner PS film and may cause surface defects [33], which  
96 are necessary for blister formation [23].

97 A HOPG (10×10 mm<sup>2</sup>, ZYH grade, SHNTI, Shanghai, China) surface was freshly cleaved as the  
98 substrate for PDMS ND nucleation. A PDMS solution with a concentration of 1/2000 (v/v) was  
99 prepared by dissolving 0.1 mL PDMS (Sylgard 184, Dow Corning, United States) into 200 mL  
100 chloroform.

## 101 **(ii) Nucleation of NBs, NDs and blisters.**

102 Generally speaking, there are two approaches to produce NBs [34]. One is the spontaneous nu-  
103 cleation by simply depositing water on hydrophobic surfaces. Surface nanocavities trap air pock-  
104 ets and lead to the nucleation of NBs [11]. The other is achieved by solubility adjustment. By ex-  
105 changing solvent or producing temperature differences, local gas supersaturation can be achieved,  
106 resulting in the nucleation of NBs. In this study, the spontaneous nucleation of NBs was adopted.  
107 The water was first kept in air for more than 10 hours for air diffusion. To avoid the possible con-  
108 tamination originating from the disposable needles of plastic syringes [22,35], a glass syringe was  
109 used to add water on the PS surface. Moreover, degassing experiments were conducted to verify  
110 the gaseous nature of the generated nanobubbles. For detailed results, please refer to SI. Blisters  
111 were formed by immersing PS sample 2 into DI water.

112 The chloroform-PDMS mixture was applied to produce PDMS NDs with controllable sizes. Since  
113 the chloroform can dissolve PS substrates, here NDs were nucleated on a HOPG substrate. First,  
114 20 microliter PDMS-chloroform solution was deposited onto the freshly cleaved HOPG substrate.  
115 After the rapid volatilization of chloroform, the substrate with remained PDMS was immersed in  
116 water by depositing a water drop (200 microliter). The PDMS NDs was then nucleated at the wa-  
117 ter/HOPG interface.

## 118 **(iii) AFM Characterization.**

119 A commercial AFM (Resolve, Bruker, U.S.A.) was used for imaging the sample in both air and  
120 DI water. Silicon cantilevers (NSC36/ALBS, MikroMasch) with the quoted tip radius of 8 nm and  
121 stiffness of 0.6 N/m from the manufacturer was used. The measured resonance frequencies of the  
122 cantilever in air and water were about 55 kHz and 16 kHz, respectively. While imaging, the driving  
123 frequencies of cantilevers were slightly lower than their resonance frequencies. A scan rate of 2  
124 Hz with a 0° scan angle was used. All experiments were performed at an ambient environment  
125 (temperature: 26 °C).

126 All sample surfaces were first scanned in air by the tapping mode AFM (TMAFM). After that, the  
127 imaging was performed in liquid. The cantilever's spring constant for scanning NBs, NDs and blis-  
128 ters are 3.2, 2.8 and 3.1 N/m, respectively, which were calibrated by a thermal noise method. The  
129 free oscillation amplitude of the cantilever at working frequency was about 200 mV. During exper-  
130 iments, the nucleated NBs, NDs and blisters were first scanned with a larger oscillation amplitude  
131 setpoint of about 195 mV (namely 98% of free amplitude) to minimize the disturbance of AFM  
132 tips to sample surfaces. To investigate the mobility of the three objects, higher scanning loads were  
133 used. After that, 98% setpoint scanning was performed again to check the consequence of high  
134 load scanning. The scan areas for NBs and NDs are  $2\ \mu\text{m} \times 2\ \mu\text{m}$ , and  $5\ \mu\text{m} \times 5\ \mu\text{m}$  for blisters.  
135 Subsequently, smaller areas (around  $600\ \text{nm} \times 600\ \text{nm}$ ) with individual NBs, NDs and blisters were  
136 selected to perform a successive scans with setpoint values varying from 98% to 60%.

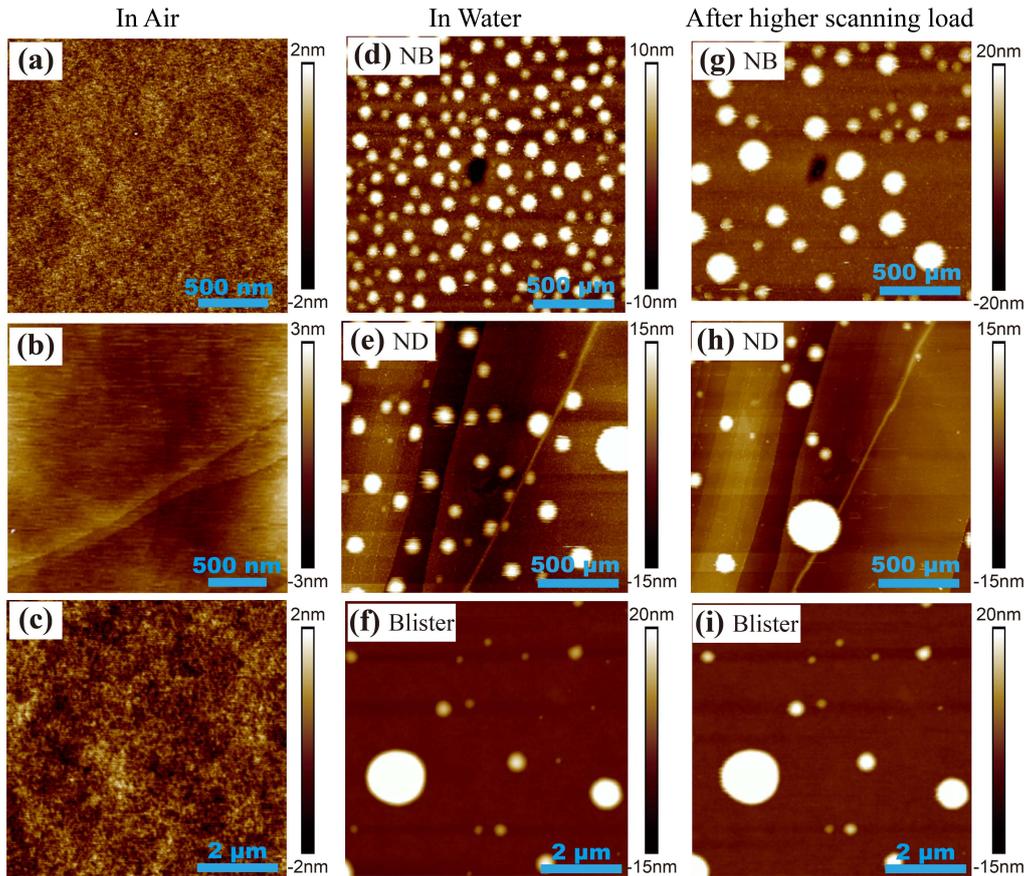
137 Force volume mode AFM was used to get several series of force-distance curves on a NB, a ND,  
138 and a blister. In the mode, the AFM cantilevers perform force-distance curve measurement in a  
139 specific area with fixed step sizes along x and y axes. At each scan position, the AFM scanner per-  
140 forms a vertical extension and retraction motion relative to the sample surface under constant driv-  
141 ing frequency and amplitude, and a force distance curve was then obtained. The ramp size for the  
142 extension and retraction motion was 200 nm and the scan rate was 1 Hz. The step sizes along both  
143 x and y axes for NBs and NDs experiments were 20 nm, within a scan area of  $500\ \text{nm} \times 500\ \text{nm}$ .  
144 The AFM images are processed with a home designed algorithm for image segmentation, contact  
145 angle measurement and volume calculation [36-39]. All calculations were performed in MATLAB  
146 2016.

## 147 **Results and Discussion**

### 148 **Morphological characterization and high load scanning**

149 The images of PS surface 1, HOPG, and PS surface 2 in air are shown in Figure 1 (a), (b) and (c),  
150 respectively. All these surfaces are smooth in air. Corresponding AFM images of NBs, NDs and

151 blisters in deionized (DI) water were depicted in Figure 1 (d), (e) and (f), respectively. To minimize  
 152 the influence of the high load scanning on the measured morphology, these images were captured  
 153 under a very low scanning load with a setpoint of 98% of free amplitude of the AFM cantilevers.  
 154 One can see that NBs and NDs have relatively uniform distribution of lateral size  $L$ , while  $L$  varies  
 155 a lot for blisters.

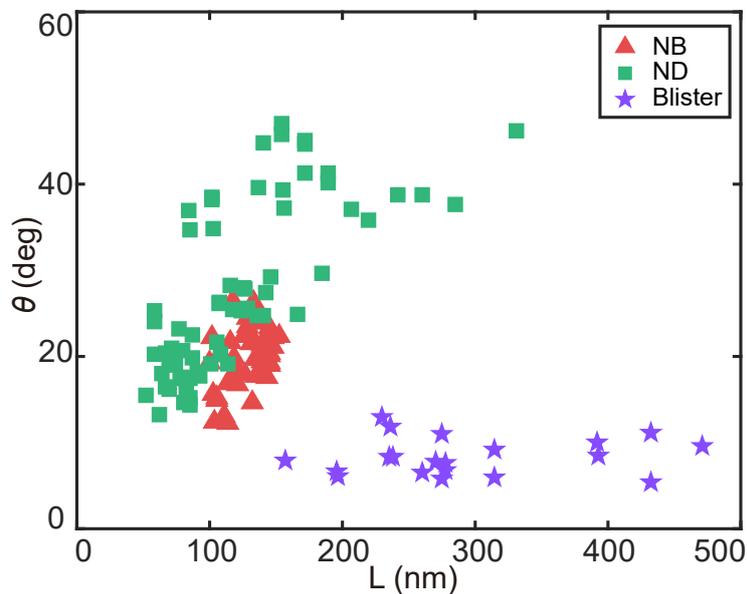


**Figure 1:** Generation of NBs, NDs, and blisters and their response to high scanning loads. (a)-(c) The PS surface 1 (a), HOPG surface (b), and PS surface 2 (c) in air. (d)-(f) The generated NBs (d), NDs (e), and blisters (f) after the surfaces were immersed in water. (g)-(i) AFM images of the NBs, NDs, and blisters after high load scans. The NBs and NDs are movable and larger NBs and NDs formed after coalescence. The blisters remain at the same positions regardless of the scanning load.

156 AFM images also reveal that the three objects exhibit different aspect ratios. Figure 2 depicts the  
 157 contact angle  $\theta$  of NBs, NDs, and blisters from the gas, oil, and solid side, respectively. It clearly  
 158 shows that  $\theta$  increases with  $L$  for both NBs and NDs. On the contrary,  $\theta$  is independent of  $L$  for

159 blisters. In addition, NBs and NDs have larger  $\theta$  than that of blisters, and  $\theta$  for NBs is slightly  
160 lower than ND.

161 The response of NBs to higher load scanning is a typical approach to study their physical prop-  
162 erties [31,38]. Here a higher scanning load (setpoint 60%) was applied for imaging. After that, a  
163 98% setpoint lower load scanning was applied to the same areas to observe the change of sample  
164 surfaces. The results are shown in Figure 1 (g-i). It shows that the NBs and NDs were moved and  
165 coalesced during the higher load scanning. As a result, larger NBs and NDs were nucleated (Figure  
166 1g and h). However, no apparent change was observed for blisters. They all remained at the same  
167 positions after the high load scanning.



**Figure 2:** Contact angles of NBs, NDs and blisters as a function of diameter. The contact angles of NBs and NDs increase with increasing lateral size, and are larger than that of blisters. The contact angle of blisters is independent of their diameter.

168 The observed coalescence of NBs is consistent with that reported somewhere else [11]. Under a  
169 higher scanning load, larger lateral force can be applied to nano-objects. For NBs, when the lateral  
170 force exceeds the pinning force origin at the solid-liquid-gas three-phase contact lines, they can be  
171 moved and hence coalesced with neighboring ones. This leads to formation of larger NBs with re-  
172 duced area density. The mechanism has been well studied in our previous studies. A term *immobil-*  
173 *ity* was proposed to associate surface nanostructures to the resistance of NBs to lateral forces [40].

174 NDs share the same mechanism of movement and coalescence under the higher scanning loads  
 175 [22]. An *et al.* reported that moving NDs is even easier than that of NBs [31]. While for blisters,  
 176 they are actually PS film wrapped water pockets [23]. They cannot be moved with a tapping mode  
 177 AFM, not to mentioned the following coalescence.

178 The above results show that NBs and NDs exhibit distinct behaviors from blisters under higher  
 179 scanning loads. Therefore, one can easily distinguish blisters from NBs and NDs. The difference  
 180 between NBs and NDs under the high scanning loads worths further investigation. Here the vol-  
 181 ume of NBs and NDs before and after coalescence were compared with a home-designed algorithm  
 182 (refer to SI for details). Interestingly, we find that the volume change before and after the coales-  
 183 cence is quite different between NBs and NDs, as shown in Table 1. The total volumes  $V_{nb}$  and  $V_{nd}$   
 184 before coalescence in Figure 1 (d) and (e) are  $4.55 \times 10^6 \text{ nm}^3$  and  $1.05 \times 10^7 \text{ nm}^3$  for NBs and NDs,  
 185 respectively. After coalescence (Figure 1 g and h), the total volumes are  $V_{nb}=1.1 \times 10^7 \text{ nm}^3$  and  
 186  $V_{nd}=1.11 \times 10^7 \text{ nm}^3$  for NBs and NDs, respectively.

**Table 1:** Volume change of NBs and NDs before and after coalescence.

	Before coalescence	After coalescence
$V_{nb} (\times 10^6 \text{ nm}^3)$	$4.55 \pm 0.29$	$11.00 \pm 0.20$
$V_{nd} (\times 10^7 \text{ nm}^3)$	$1.05 \pm 0.03$	$1.11 \pm 0.02$

187 The increment of NB volume is highly repeatable among different experiments (see the repeated  
 188 experiment in SI). We believe that this is because of the reduced inner pressure with increasing NB  
 189 size. According to Young-Laplace equation,

$$190 \quad p_g = p_0 + \frac{2\gamma_L}{R}, \quad (1)$$

191 where  $p_0$  is the ambient pressure,  $\gamma_L$  is surface tension of water, and  $R$  is the radius of NBs. As-  
 192 suming that the NBs are equilibrated with each other and the gas concentration in the edge of  
 193 bubbles is constant, the inner pressure  $p_g$  rapidly reduces with increasing  $R$ . This has two con-

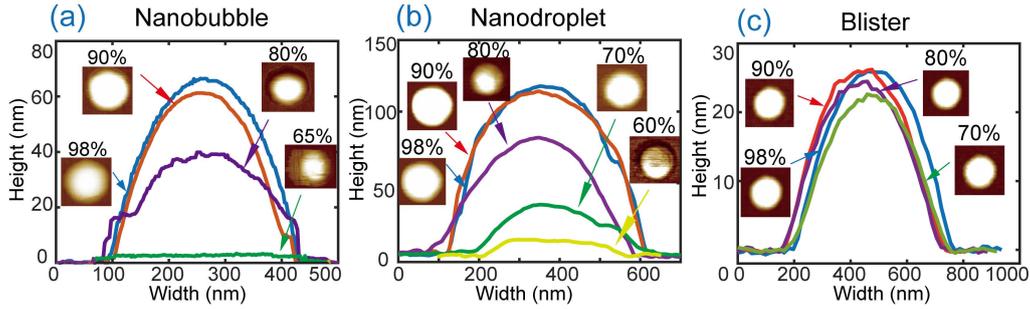
194 sequences. First, based on the ideal gas law, the gas molecules originally trapped in NBs before  
195 coalescence will occupy more space when  $p_g$  is reduced. This leads to increased bubble volumes.  
196 Moreover, the reduced inner pressure  $p_g$  enhances the diffusion of the dissolved gas in water into  
197 NBs. This further increases NB volume. Regarding the NDs, as expected, the volumes are almost  
198 the same before and after the coalescence. The results indicate that the content in the spherical ob-  
199 jects on PS surface 1 are indeed gas.

200 After the 60% setpoint high load scanning, the successive AFM scans with setpoints from 98% to  
201 60% of the free oscillation amplitude were conducted. The morphology of NBs and NDs changed  
202 significantly with decreasing setpoints, *i. e.* imaging forces, as shown in figure 3(a) and (b). The  
203 cross section profiles of NB and ND height images decreases with increasing scanning forces. The  
204 cross section profiles obtained with setpoints larger than 80% still exhibit circular arc shape. No-  
205 tably, the NB become almost flat at the lowest setpoint of 65%, while the ND appeared as a som-  
206brero (a spherical cap sitting on a flat molecular layer) at setpoints of 70% and 60%. Comparing  
207 with the NB, the ND has larger resistance to vertical loads and can still maintain its circular arc  
208 shape for setpoints down to 60%. This is because the gaseous NBs are softer than NDs, thus easier  
209 to be deformed by AFM tips [31,34,41]. Particularly, the remained molecular layer of ND under  
210 larger vertical loads have been observed in several previous studies [31,42], which is consistent  
211 with our results, although the underlying mechanism is still not well understood. Unsurprisingly,  
212 the height of blister in figure 3 (c) only slightly decreases in height with increasing scanning forces,  
213 exhibiting a solid-like response to scanning loads.

## 214 **Force spectroscopy**

215 To distinguish one from the other in between NBs and NDs, the force spectroscopy characterization  
216 was conducted through tip-sample interaction measurement using the AFM. By tracking the ver-  
217 tical deflection signal of a AFM cantilever as it approaches and retracts from sample surfaces, the  
218 mechanical signatures of the three objects were revealed.

219 Figure 4 (a)-(c) depict three typical force-distance curves on a NB, a ND, and a blister, respectively.

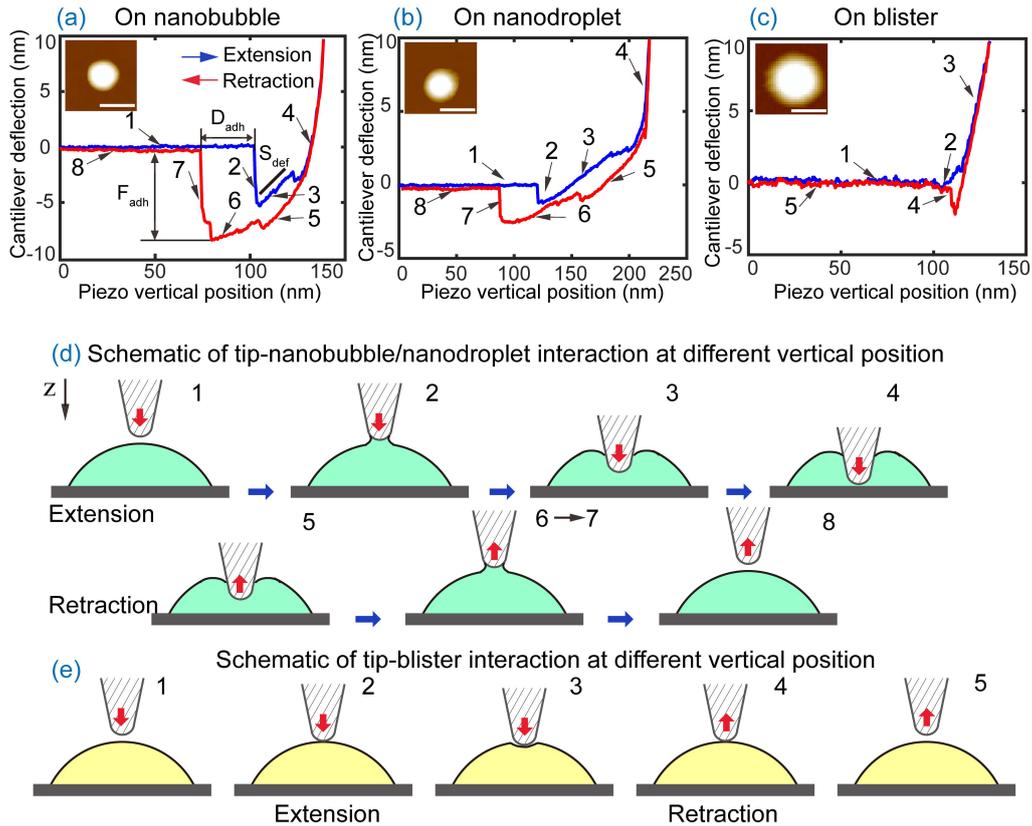


**Figure 3:** Morphology change of the three different surface nano-objects under various scanning loads. The cross section profiles of a NB (a), a ND (b) and a blister (c) with increasing scanning loads of setpoints from 98% to 60%. The height profiles of the NB and ND both reduce with the increasing imaging force, while the height for the blister does not change much.

220 The insets are corresponding AFM height images of the three nano-objects. Here a NB and a ND  
 221 with a similar footprint diameter were produced through coalescence to minimize the influence of  
 222 their sizes on the characterization. Figure 4(d) and (e) are schematics of the different stages of tip-  
 223 NB/ND and tip-blister interaction, respectively. Note here that the downward motion direction of  
 224 the  $z$  piezo stage is defined as the positive direction.

225 As shown in section 1 of figure 4 (a), when the tip is far away from the substrate (in liquid), there is  
 226 no interaction between them and thus the deflection remains constant at zero (figure 4 (d)-1). Once  
 227 the AFM tip contacts the NB, the tip is rapidly drawn into the NB due to the capillary force (figure  
 228 4 (d)-2). Therefore, an obvious snap-in is observed in section 2. This is consistent with the previ-  
 229 ous studies [40,43-45]. In section 3, the deflection linearly increases with piezo vertical position  $Z$ .  
 230 That is because that the perimeter of the three-phase contact line increases with insertion depth of  
 231 the cantilever tip into the NB. As a result, the force applied to the tip from the three-phase contact  
 232 line also show linear dependence on  $z$ , as illustrated in figure 4(d)-3.

233 When the tip further goes down, it touches the solid sample surface (figure 4 (d)-4). This hard con-  
 234 tact results in a rapid linear increase of cantilever deflection with increasing piezo vertical position  
 235  $z$  in section 4. After the AFM  $z$  scanner goes to the minimum vertical position, it reverses its mo-  
 236 tion and starts the retraction motion. It starts from the rapid linear decrease of the deflection signal.  
 237 After that, it enters the linear reduction part (section 5, figure 4 (d)-5), until the tip snaps out of the  
 238 NB (section 6, figure 4 (d)-6). Finally, the deflection become zero in section 7 (figure 4 (d)-7).



**Figure 4:** Force spectroscopy measurement through force-distance curves. (a-c) Typical force-distance curves obtained on a NB (a), a ND (b) and a blister (c). (d) The schematic of tip-NB/ND interaction. The numbers 1-8 correspond to the different stages labeled on the force distance curves in (a) and (b). (e) The schematic of tip-blister interaction, indicating the five stages along with the force curve.

239 Notably, the AFM tip used here is silicon, which is hydrophilic. Previous studies demonstrate that  
 240 the hydrophobicity of AFM tips can significantly influence tip-bubble interactions and hence the  
 241 bubble imaging [46,47]. For hydrophobic tips the bubble interface may jump towards the tips. This  
 242 results in a large snap-in of the force curves and cause larger bubble deformation. In contrast, scan-  
 243 ning with a hydrophilic tip provides reliable tip-bubble interaction curves as well as less distorted  
 244 bubble images. In addition, the linearity of the tip-bubble interaction region also depends on the  
 245 hydrophobicity of the tips [45]. With a hydrophilic tip, the tip-bubble interaction shows better lin-  
 246 earity and exhibits a lower slope value.

247 The tip-ND interaction is very similar to the tip-NB interaction (figure 4 (b)). However, tip-blister

248 interaction exhibits distinct behaviors from that of tip-NB/ND interactions, as shown in figure 4 (c).

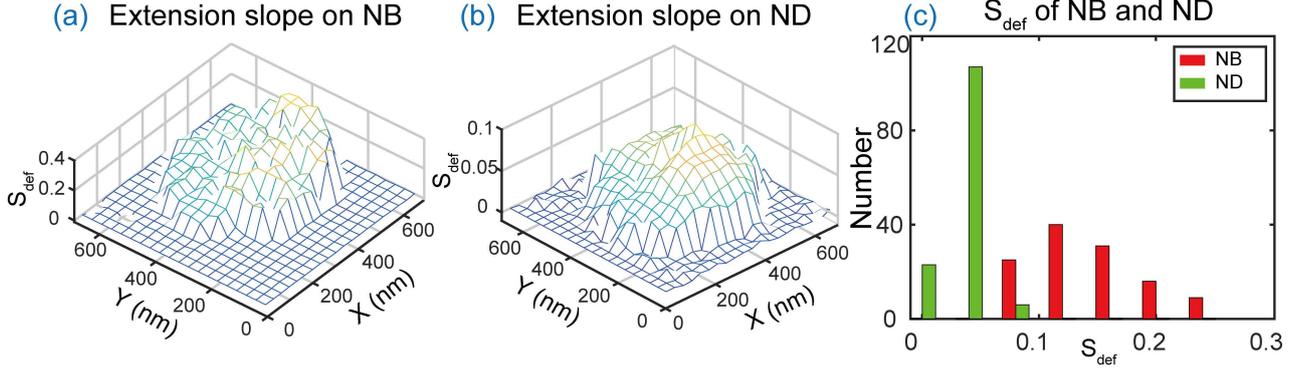
249 The force curves measured on the blisters is close to that measured on the bare PS substrate. The  
250 schematics in figure 4 (e) illustrate the different stages of tip-blisters interaction in figure 4 (c). One  
251 can see that blisters more behave as solid objects in tip-sample interaction measurement.

252 With the above force spectroscopy measurement, it is easy to distinguish NBs from blisters. How-  
253 ever, force-distance curves on the soft objects of NBs and NDs are very close to each other. In or-  
254 der to further investigate the difference between NBs and NDs, we extract several key indicators  
255 from force distance curves.

256 As shown in Figure 4 (a), the prefactor  $S_{def}$  in the linear interaction region of section 3, the max-  
257 imum adhesion force  $F_{adh}$  and the distance  $D_{adh}$  can be obtained for individual force distance  
258 curves. By performing the force volume mode measurement, a series of force-distance curves at  
259 specific positions of NBs/NDs were acquired. For each force-distance curve, the three parameters  
260 were manually extracted. The prefactor  $S_{def}$  in section 3 is the derivative of vertical forces to verti-  
261 cal piezo distance  $z$ . According to Ref [45],  $S_{def} \propto \gamma_L$ , where  $\gamma_L$  is the liquid/gas surface tension  
262 of water. Thus the prefactor  $S_{def}$  indicates the surface tension of the object which the tip penetrates  
263 into. The surface tension of water/air and water/PDMS interfaces are  $72 \text{ mNm}^{-1}$  and  $40 \text{ mNm}^{-1}$ ,  
264 respectively. Figure 5 (a) and (b) are the constructed maps of  $S_{def}$  on NBs and NDs, respectively.  
265 Figure 5 (c) shows the histograms of  $S_{def}$  on the NB and the ND. One can see that  $S_{def}$  on the NB  
266 is much larger than that on the ND, which is consistent with the derived mathematic model in Ref  
267 [45].

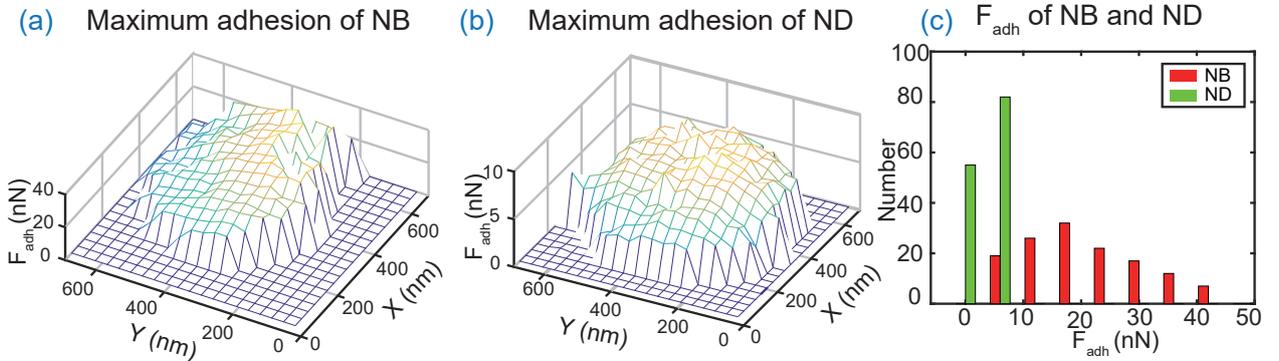
268 The value  $F_{adh}$  is the adhesion force between the AFM tip and the NB/ND at the moment of snap-  
269 out. Since the radius of the NB is about 1000 times larger than that of the AFM tip, the NB can  
270 be approximated as a planar surface when modeling the tip-NB interaction. By assuming that the  
271 tip was in contact with a continuous film, the attractive force  $F_{adh}$  due to the meniscus bridge for a  
272 sphere in contact with a plane surface can be expressed as [45]

$$273 \quad F_{adh} = 2\pi R\gamma_L(1 + \cos\theta), \quad (2)$$



**Figure 5:** Comparison of the prefactor  $S_{def}$  in the linear interaction region of force-distance curves between a NB and a ND. The constructed maps of  $S_{def}$  for the NB (a) and the ND (b). (c) The histograms of  $S_{def}$  on the NB and the ND. The NB exhibits much larger value of  $S_{def}$ .

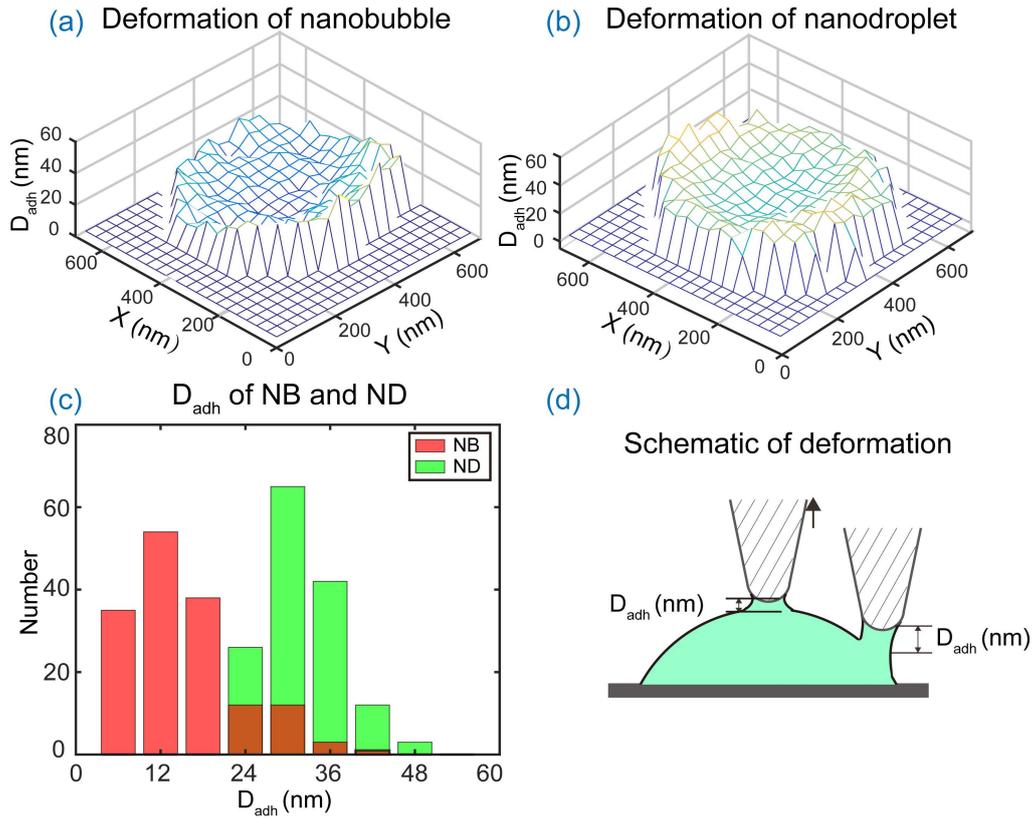
274 where  $R$  is the radius of the sphere (tip radius here),  $\gamma_L$  is the water/gas (NB case) or water/PDMS  
 275 (ND case) surface tension, and  $\theta$  is the contact angle of AFM cantilever on the NB or ND. The  
 276 constructed maps of  $F_{adh}$  on the NB and the ND are shown in Figure 6 (a) and (b), respectively.  
 277 We find that the NB have much higher value of  $F_{adh}$  comparing to the ND. This is consistent with  
 278 the reported surface tension values at the three phase contact lines of the NB and ND.



**Figure 6:** Comparison of the tip-sample adhesion force between the NB and the ND. The constructed maps of the adhesion force  $F_{adh}$  for the NB (a) and the ND (b). (c) The histograms of the adhesion force  $F_{adh}$  on the NB and the ND. The NB has much larger  $F_{adh}$ .

279 The parameter  $D_{adh}$  extracted from the force distance curves is the length of pulled capillary  
 280 bridge, namely, the deformation of the NB/ND when the tip detaches from them. In most cases,  
 281 the capillary bridge can be pulled out for a limited distance before the AFM tip eventually detaches.

282 As shown in Figure 7, the  $D_{adh}$  of the NB is smaller than that of the ND. This result indicates that  
 283 the ND has a much larger capillary bridge than that of the NB. The medium value of  $D_{adh}$  on the  
 284 ND is about 36 nm, which is about three times of the value measured on the NB. This is because  
 285 that the droplet of PDMS is non-Newtonian liquid. That can stretch longer, since the stretching  
 286 modifies the stress balance anisotropically.



**Figure 7:** Comparison of deformation between the NB and the ND. The constructed maps of the deformation  $D_{adh}$  for the NB (a) and the ND (b). (c) The comparison of the histograms of  $D_{adh}$  on the NB and ND. (d) Schematic of the deformation of NB/ND by the AFM tip. Due to the higher viscosity value of PDMS than water,  $D_{adh}$  on the ND is much larger than that on the NB.

287 From the above results, one can see that the three nano-objects are distinguishable by using one ap-  
 288 proach or the combination of several approaches through morphological characterization and force  
 289 spectroscopy measurement. The distinguishability of one object to the other two objects is summa-  
 290 rized in table 2. One can see that blisters have the highest distinguishability among the three nano-  
 291 objects. They can be distinguished from the other two kinds of nano-objects, no matter based on

CA measurement, high load scanning, or force spectroscopy measurement. The challenge remains on the distinction of NBs and NDs from each other. The major reason is that both NBs and NDs are all soft in nature and have either liquid/gas or liquid/liquid interfaces. As a result, they show similar response to high load scanning, exhibiting high mobility and deformability. Fortunately, they are distinguishable through volume measurement in coalescence experiment. NBs show significantly increased volume after coalescence, while the volume of NDs remains the same before and after coalescence. Regarding force spectroscopy measurement, NBs and NDs can be distinguished through quantitative evaluation, based on the fact that the surface tension value in the two cases is different.

**Table 2:** The distinguishability of the nano-objects.

	CA	Response to high load			Force spectroscopy		
		Mobility	Deformability	Volume	$S_{\text{def}}$	$F_{\text{adh}}$	$D_{\text{adh}}$
Nanobubble	✗	✗	✗	✓	✓	✓	✓
Nanodroplet	✗	✗	✗	✓	✓	✓	✓
Blister	✓	✓	✓	—	✓	✓	✓

## Conclusion

In this study, the distinction among surface NBs, NDs and blisters was systematically investigated with an AFM through three approaches: morphological characterization, high load scanning, and force spectroscopy measurement. The results show that blisters can be easily distinguished from the other two nano-objects. They have the lowest contact angle values and are not movable and deformable under high scanning loads. The force spectroscopy measurement on the blisters is close to that obtained on solid surface.

The NBs and NDs have similar contact angles. They are all movable and deformable, and exhibit similar response at force-distance curves. However, they can be well distinguished by volume and

310 force spectroscopy measurement. The volume of NBs significantly increased after coalescence,  
311 while it remains the same for NDs. This is because of the reduced inner pressure and gas diffusion  
312 from liquid to the coalesced NBs. In the force spectroscopy measurement, three parameters, the  
313 prefactor  $S_{def}$  of linear tip-NB/ND interaction region, the adhesion force  $F_{adh}$ , and the deforma-  
314 tion  $D_{adh}$  are extracted. The results show that NBs have larger  $S_{def}$ ,  $F_{adh}$ , which agrees well with  
315 the fact that the water/air surface tension is higher than water/oil. On the contrary,  $D_{adh}$  on NBs is  
316 much smaller than that on NDs, due to the lower viscosity of water. We believe that this work pro-  
317 vides a useful approach to identify NBs, NDs and blisters and is particularly useful in surface NB  
318 studies.

## 319 **Supporting Information**

320 Supporting Information File 1:

321 File Name: S1.pdf

322 File Format: PDF

323 Title: Supplementary materials

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